

Semi - Lagrange Notes

These notes are intended for use by coders.

The equations are written using non - standard multi - character symbols which appear in the computer code. These equations have first been derived using vector notation and subsequently expanded for detailed coding.

Stable extrapolation two time level scheme SETTLS

$$\vec{R}_A^{t+\Delta t} \approx \vec{R}_D^t + \Delta t \cdot \left[\frac{d\vec{R}}{dt} \right]_D^t + \frac{\Delta t^2}{2} \cdot \left[\frac{d^2\vec{R}}{dt^2} \right]_M^{t+\frac{\Delta t}{2}}$$



$$\vec{R}_A^{t+\Delta t} \approx \vec{R}_D^t + \Delta t \cdot \vec{V}_D^t + \frac{\Delta t^2}{2} \cdot \left[\frac{d\vec{V}}{dt} \right]_M^{t+\frac{\Delta t}{2}}$$

$$\left[\frac{d\vec{V}}{dt} \right]_M^{t+\frac{\Delta t}{2}} = ?$$

$$\left[\frac{d\vec{V}}{dt} \right]_M^{t+\frac{\Delta t}{2}} \approx \left[\frac{d\vec{V}}{dt} \right]_M^{t-\frac{\Delta t}{2}} = \frac{\vec{V}_A^t - \vec{V}_D^{t-\Delta t}}{\Delta t}$$

$$\vec{R}_A^{t+\Delta t} = \vec{R}_D^t + \frac{\Delta t}{2} \cdot ([2\vec{V}_D^t - \vec{V}_D^{t-\Delta t}]_D + \vec{V}_A^t)$$

Two Time Level Semi-Lagrangian Semi-implicit With SETTLS

$$\begin{aligned}
 \frac{dX}{dt} &= R = N + L \\
 \text{explicit} \quad \frac{(X_{t+\Delta t}^A - X_t^D)}{\Delta t} &= N_{t+\frac{\Delta t}{2}}^M + L_{t+\frac{\Delta t}{2}}^M \\
 \text{semi-implicit} \quad \frac{(X_{t+\Delta t}^A - X_t^D)}{\Delta t} &= N_{t+\frac{\Delta t}{2}}^M + \frac{1}{2}(L_{t+\Delta t}^A + L_t^D) \\
 &= N_{t+\frac{\Delta t}{2}}^M + \left(L_{t+\frac{\Delta t}{2}}^M - L_{t+\frac{\Delta t}{2}}^M \right) + \frac{1}{2}(L_{t+\Delta t}^A + L_t^D) \\
 &= N_{t+\frac{\Delta t}{2}}^M + L_{t+\frac{\Delta t}{2}}^M + \frac{1}{2} \left[L_{t+\Delta t}^A - 2L_{t+\frac{\Delta t}{2}}^M + L_t^D \right] \\
 &= \text{explicit} + \text{correction} \\
 \text{correction} &= \frac{1}{2} \left[L_{t+\Delta t}^A - 2L_{t+\frac{\Delta t}{2}}^M + L_t^D \right] = \frac{1}{2} \Delta_{tt} L
 \end{aligned}$$

The correction can be viewed as semi-Lagrange Laplacian.

The correction can also be weighted with a factor β .

$$\frac{dX}{dt} = (N + L)_{t+\frac{\Delta t}{2}}^M + \frac{1}{2} \beta \Delta_{tt} L$$

The terms of the form $F_{t+\frac{\Delta t}{2}}^M$ are approximated :

$$F_{t+\frac{\Delta t}{2}}^M \approx \frac{1}{2} \left(F_t^A + (2F_t - F_{t-\Delta t})^D \right) \quad (\text{ECMWF Settls approximation})$$

Let $\dot{X} = N+L$

$$X_{t+\Delta t}^A = X_t^D + \frac{\Delta t}{2} \left(\dot{X}_t^A + (2\dot{X}_t - \dot{X}_{t-\Delta t})^D \right) + \frac{1}{2} \beta \Delta t \Delta_{tt} L$$

$$X_{t+\Delta t}^A = \frac{\Delta t}{2} \dot{X}_t^A + \Delta t \left[\frac{X_t}{\Delta t} + \dot{X}_t - \frac{1}{2} \dot{X}_{t-\Delta t} \right]^D + \frac{1}{2} \beta \Delta t \Delta_{tt} L$$

expand :

$$\begin{aligned} \Delta_{tt} L &= \left[L_{t+\Delta t}^A - 2L_{t+\frac{\Delta t}{2}}^M + L_t^D \right] \\ &= L_{t+\Delta t}^A - 2 \left\{ \frac{1}{2} \left[L_t^A + (2L_t - L_{t-\Delta t})^D \right] \right\} + L_t^D \end{aligned}$$

$$\Delta_{tt} L = L_{t+\Delta t}^A - L_t^A + (L_{t-\Delta t} - L_t)^D$$

The surface pressure equation.

We need an expression for $\frac{dp_s}{dt}$.

start with continuity :

$$\frac{\partial}{\partial t} \frac{\partial p}{\partial \eta} + \nabla_3 \cdot \frac{\partial p}{\partial \eta} \bar{\mathbf{V}}_3 = 0, \text{ expand and regroup:} \quad 1$$

$$\frac{d}{dt} \frac{\partial p}{\partial \eta} = - \frac{\partial p}{\partial \eta} \left(D + \frac{\partial \dot{\eta}}{\partial \eta} \right), \quad D = \nabla \cdot \bar{\mathbf{v}}_H \quad 2$$

now also express $\frac{d}{dt} \frac{\partial p}{\partial \eta}$ using $p = A(\eta) + B(\eta)p_s$

then

$$\frac{d}{dt} \frac{\partial p}{\partial \eta} =$$

$$\frac{\partial}{\partial t} \frac{\partial}{\partial \eta} (A + Bp_s) + \bar{\mathbf{V}} \cdot \nabla \frac{\partial}{\partial \eta} (A + Bp_s) + \dot{\eta} \frac{\partial}{\partial \eta} \frac{\partial p}{\partial \eta}$$

$$= \frac{\partial B}{\partial \eta} \frac{\partial p_s}{\partial t} + \bar{\mathbf{V}} \cdot \nabla \frac{\partial B}{\partial \eta} p_s + \dot{\eta} \frac{\partial}{\partial \eta} \frac{\partial p}{\partial \eta}$$

$$= \frac{\partial B}{\partial \eta} \left(\frac{\partial p_s}{\partial t} + \vec{V} \cdot \nabla p_s + \dot{\eta} \frac{\partial p_s}{\partial \eta} \right) + \dot{\eta} \frac{\partial}{\partial \eta} \frac{\partial p}{\partial \eta}$$

$$\frac{d}{dt} \frac{\partial p}{\partial \eta} = \frac{\partial B}{\partial \eta} \frac{d p_s}{dt} + \dot{\eta} \frac{\partial}{\partial \eta} \frac{\partial p}{\partial \eta}$$

Equate the two expressions for $\frac{d}{dt} \frac{\partial p}{\partial \eta}$

$$\frac{\partial B}{\partial \eta} \frac{d p_s}{dt} + \dot{\eta} \frac{\partial}{\partial \eta} \frac{\partial p}{\partial \eta} = - \frac{\partial p}{\partial \eta} \left(D + \frac{\partial \dot{\eta}}{\partial \eta} \right) \quad 3$$

$$\frac{\partial B}{\partial \eta} \frac{d p_s}{dt} + \frac{\partial p}{\partial \eta} D + \frac{\partial}{\partial \eta} \left(\dot{\eta} \frac{\partial p}{\partial \eta} \right) = 0$$

Vertical discretization yields :

$$\Delta B_k \frac{d p_s}{dt} + \Delta p_k D_k + \left(\dot{\eta} \frac{\partial p}{\partial \eta} \right)_{k+\frac{1}{2}} - \left(\dot{\eta} \frac{\partial p}{\partial \eta} \right)_{k-\frac{1}{2}} = 0 \quad 4$$

for $k = 1, \dots, K$. Recall that $\sum_{k=1}^K \Delta B_k = 1$

Divide eq. 4 by p_s to get $\frac{d}{dt} \ln p_s$

$$\Delta B_k \frac{d \ln p_s}{dt} + \frac{1}{p_s} \left[\Delta p_k D_k + \left(\dot{\eta} \frac{\partial p}{\partial \eta} \right)_{k+\frac{1}{2}} - \left(\dot{\eta} \frac{\partial p}{\partial \eta} \right)_{k-\frac{1}{2}} \right] = 0 \quad 5$$

From the vertical finite difference :

$$\left(\dot{\eta} \frac{\partial p}{\partial \eta} \right)_{k + \frac{1}{2}} = - p_s \left\{ B_{k + \frac{1}{2}} \frac{\partial q}{\partial t} + \sum_{j=1}^k \left[\frac{1}{p_s} D_j \Delta p_j + C_j \Delta B_j \right] \right\} \quad 6$$

where $q = \ln p_s$ $C_j = \vec{V}_j \cdot \nabla \ln p_s$

Substitute (6) in 5 :

$$\Delta B_k \frac{dq}{dt} + \frac{1}{p_s} \Delta p_k D_k - \left[B_{k + \frac{1}{2}} \frac{\partial q}{\partial t} + \sum_{j=1}^k \right] + \left[B_{k - \frac{1}{2}} \frac{\partial q}{\partial t} + \sum_{j=1}^{k-1} \right] = 0$$

which simplifies to :

$$\Delta B_k \frac{dq}{dt} + \frac{1}{p_s} \Delta p_k D_k - \left(B_{k+\frac{1}{2}} - B_{k-\frac{1}{2}} \right) \frac{\partial q}{\partial t} - \frac{1}{p_s} D_k \Delta p_k - C_k \Delta B_k = 0$$

$$\Delta B_k \frac{dq}{dt} - \Delta B_k \left(\frac{\partial q}{\partial t} + C_k \right) = 0 \quad 7$$

$$\text{Here } \frac{\partial q}{\partial t} = - \sum_{k=1}^K \left(\frac{1}{p_s} D_k \Delta p_k + C_k \Delta B_k \right) \quad 8$$

as required by the vertical finite differencing.

We now discretize in time while adding the correction term at every level. The final prediction eq. for the surface pressure is :

$$q_{t+\Delta t}^A = \sum_{k=1}^K \Delta B_k \left\{ q_t^D + \Delta t [Z(k)]_{t+\frac{\Delta t}{2}}^M - \Delta t \frac{\beta}{2} \Delta_{tt} L \right\} \quad 9$$

Here $Z(k) = \frac{\partial q}{\partial t} + C_k$, $\frac{\partial q}{\partial t}$ is given by eq. 8

$$L = \frac{1}{p_s^r} \sum_{j=1}^{levs} (\Delta p_j^r D_j)$$

Evaluation of the RHS of eq. 9 :

The $Z(k)$ term :

$$\Delta t [Z(k)]_{t+\frac{\Delta t}{2}}^M = \frac{\Delta t}{2} [Z(k)_t^A + (2Z(k)_t - Z(k)_{t-\Delta t})^D]$$

$$= \frac{\Delta t}{2} Z(k)_t^A + \Delta t \left[Z(k)_t - \frac{1}{2} Z(k)_{t-\Delta t} \right]^D$$

The correction term :

$$-\frac{\beta}{2} \Delta t \Delta_{tt} \left[\frac{1}{p_s^r} \sum_{j=1}^K \Delta p_j^r D_j \right] = -\Delta t \Delta_{tt} \delta$$

$$\text{where } \delta = \frac{\beta}{2 p_s^r} \sum_{j=1}^K \Delta p_j^r D_j = S \cdot \tilde{D}$$

$$S = \frac{\beta}{2 p_s^r} (\Delta p_1^r \quad \dots \quad \Delta p_K^r), \quad D = (D_1 \quad \dots \quad D_K)$$

Note that after the vertical summation δ is a 2 dimensional field.

$$\text{Since : } \Delta_{tt} \delta = \delta_{t+\Delta t}^A - 2 \delta_{t+\frac{\Delta t}{2}}^M + \delta_t^D$$

$$\text{and : } \delta_{t+\frac{\Delta t}{2}}^M = \frac{1}{2} \left[\delta_t^A + (2\delta_t - \delta_{t-\Delta t})^D \right]$$

$$\Delta_{tt} \delta = \delta_{t+\Delta t}^A - 2 \cdot \frac{1}{2} \left[\delta_t^A + (2\delta_t - \delta_{t-\Delta t})^D \right] + \delta_t^D$$

$$\Delta_{tt} \delta = \delta_{t+\Delta t}^A - \delta_t^A + (\delta_{t-\Delta t} - \delta_t)^D$$

The final form of eq. 9 is therefore :

$$q_{t+\Delta t}^A = \sum_{k=1}^K \Delta B_k \left\{ q_t^D + \frac{\Delta t}{2} Z(k)_t^A + \Delta t \left[Z(k)_t - \frac{1}{2} Z(k)_{t-\Delta t} \right]^D \right. \\ \left. - \Delta t \left(\delta_{t+\Delta t}^A - \delta_t^A + (\delta_{t-\Delta t} - \delta_t)^D \right) \right\}$$

Grouping arrival and departure point contributions :

$$q_{n+1}^A = \sum_{k=1}^K \Delta B_k \left\{ \frac{\Delta t}{2} Z(k)_t^A - \Delta t \delta_{t+\Delta t}^A + \Delta t \delta_t^A \right\} \\ + \sum_{k=1}^K \Delta B_k \left\{ q_t + \Delta t \left[Z(k)_t - \frac{1}{2} Z(k)_{t-\Delta t} + \delta_t - \delta_{t-\Delta t} \right]^D \right\}$$

For the semi-implicit solution we write, noting that the vertical sum at arrival points collapses for δ_t^A and $\delta_{t+\Delta t}^A$:

$$Z = \Delta t \delta_t^A + \sum_{k=1}^K \frac{\Delta t}{2} \Delta B_k Z(k)_t +$$

$$\sum_{k=1}^K \Delta B_k \left\{ q_t + \Delta t \left[Z(k)_t - \frac{1}{2} Z(k)_{t-\Delta t} + \delta_t - \delta_{t-\Delta t} \right] \right\}^D$$

$$q_{t+\Delta t}^A = Z - \Delta t \delta_{t+\Delta t}^A = Z - \frac{\beta \Delta t}{2} \frac{1}{p_s^r} \sum_{j=1}^K \Delta p_j^r (D_j)_{t+\Delta t}^A$$

$$\frac{U_{t+\Delta t}^A - U_t^D}{\Delta t} = \left\{ fV - \frac{1}{a} \left[\frac{\partial \phi}{\partial \lambda} + R_d T_v \frac{\partial}{\partial \lambda} \ln p \right] \right\}_{t+\frac{\Delta t}{2}}^M$$

$$- \frac{\beta}{2a} \Delta_{tt} \left[Y \frac{\partial T}{\partial \lambda} + R_d T_r \frac{\partial}{\partial \lambda} \ln p_s \right]$$

$$UN = fV - \frac{1}{a} \left[\frac{\partial \phi}{\partial \lambda} + R_d T_v \frac{\partial}{\partial \lambda} \ln p \right]$$

$$UL = \frac{1}{a} \left[Y \frac{\partial T}{\partial \lambda} + R_d T_r \frac{\partial \ln p_s}{\partial \lambda} \right]$$

$$\Delta_{tt} f = \left(f_{t+\Delta t}^A - 2f_{t+\frac{\Delta t}{2}}^M + f_t^D \right)$$

$$\frac{U_{t+\Delta t}^A - U_t^D}{\Delta t} = UN_{t+\frac{\Delta t}{2}}^M - \frac{\beta}{2} \left[UL_{t+\Delta t}^A - 2UL_{t+\frac{\Delta t}{2}}^M + UL_t^D \right]$$

$$= (UN + \beta UL)_{t+\frac{\Delta t}{2}}^M - \frac{\beta}{2} \left[UL_{t+\Delta t}^A + UL_t^D \right] =$$

$$UNL = UN + \beta UL$$

$$= (UNL)_{t+\frac{\Delta t}{2}}^M - \frac{\beta}{2} [UL_{t+\Delta t}^A + UL_t^D]$$

$$= \frac{1}{2} [UNL_t^A + (2UNL_t - UNL_{t-\Delta t})^D] - \frac{\beta}{2} UL_{t+\Delta t}^A - \frac{\beta}{2} UL_t^D$$

$$= \frac{1}{2} UNL_t^A + \left[\left(UNL_t - \frac{1}{2} UNL_{t-\Delta t} \right) - \frac{\beta}{2} UL_t \right]^D - \frac{\beta}{2} UL_{t+\Delta t}^A$$

save UNL_t to be used in next step as $UNL_{t-\Delta t}$

$$U_{t+\Delta t}^A + \frac{\Delta t \beta}{2} UL_{t+\Delta t}^A =$$

$$\frac{\Delta t}{2} UNL_t^A + \left(U_t + \Delta t \left(UNL_t - \frac{1}{2} UNL_{t-\Delta t} - \frac{\beta}{2} UL_t \right) \right)^D = \tilde{U}$$

$$\frac{V_{t+\Delta t}^A - V_t^D}{\Delta t} = \left\{ -fU - \frac{\cos(\theta)}{a} \left[\frac{\partial \varphi}{\partial \theta} + R_d T_v \frac{\partial}{\partial \theta} \ln p \right] \right\}_{t+\frac{\Delta t}{2}}^M$$

$$- \frac{\beta \cos(\theta)}{2a} \Delta_{tt} \left[Y \frac{\partial T}{\partial \theta} + R_d T_r \frac{\partial}{\partial \theta} \ln p_s \right]$$

$$VN = -fU - \frac{\cos(\theta)}{a} \left[\frac{\partial \varphi}{\partial \theta} + R_d T_v \frac{\partial}{\partial \theta} \ln p \right]$$

$$VL = \frac{\cos(\theta)}{a} \left[Y \frac{\partial T}{\partial \theta} + R_d T_r \frac{\partial \ln p_s}{\partial \theta} \right]$$

$$\Delta_{tt} f = \left(f_{t+\Delta t}^A - 2f_{t+\frac{\Delta t}{2}}^M + f_t^D \right)$$

$$\frac{V_{t+\Delta t}^A - V_t^D}{\Delta t} = VN_{t+\frac{\Delta t}{2}}^M - \frac{\beta}{2} \left[VL_{t+\Delta t}^A - 2VL_{t+\frac{\Delta t}{2}}^M + VL_t^D \right]$$

$$= (VN + \beta VL)_{t+\frac{\Delta t}{2}}^M - \frac{\beta}{2} \left[VL_{t+\Delta t}^A + VL_t^D \right] =$$

$$VNL = VN + \beta VL$$

$$= (VNL)_{t+\frac{\Delta t}{2}}^M - \frac{\beta}{2} [VL_{t+\Delta t}^A + VL_t^D]$$

$$= \frac{1}{2} [VNL_t^A + (2VNL_t - VNL_{t-\Delta t})^D] - \frac{\beta}{2} VL_{t+\Delta t}^A - \frac{\beta}{2} VL_t^D$$

$$= \frac{1}{2} VNL_t^A + \left[\left(VNL_t - \frac{1}{2} VNL_{t-\Delta t} \right) - \frac{\beta}{2} VL_t \right]^D - \frac{\beta}{2} VL_{t+\Delta t}^A$$

save VNL_t to be used in next step as $VNL_{t-\Delta t}$

$$V_{t+\Delta t}^A + \frac{\Delta t \beta}{2} VL_{t+\Delta t}^A =$$

$$\frac{\Delta t}{2} VNL_t^A + \left(V_t + \Delta t \left(VNL_t - \frac{1}{2} VNL_{t-\Delta t} - \frac{\beta}{2} VL_t \right) \right)^D = \tilde{v}$$

$$\text{Let } TN = \frac{\kappa T_v \omega}{p} \left(\frac{1 + \varepsilon q}{1 + (\delta - 1)q} \right), \quad TL = \tau D$$

where τ is the matrix derived in the Eulerian model

$$\text{The thermodynamic eq: } \frac{dT_v}{dt} = TN - \frac{\beta}{2} \Delta t \tau D$$

$$\frac{T_{t+\Delta t}^A - T_t^D}{\Delta t} = TN_{t+\frac{\Delta t}{2}}^M - \frac{\beta}{2} \left(TL_{t+\Delta t}^A - 2TL_{t+\frac{\Delta t}{2}}^M + TL_t^D \right)$$

$$= (TN + \beta TL)_{t+\frac{\Delta t}{2}}^M - \frac{\beta}{2} (TL_{t+\Delta t}^A + TL_t^D) =$$

$$\text{Let } TNL = TN + \beta TL \quad \text{initialize } TNL(t - \Delta t)$$

$$= \frac{1}{2} \left[TNL_t^A + (2TNL_t - TNL_{t-\Delta t})^D \right] - \frac{\beta}{2} (TL_{t+\Delta t}^A + TL_t^D)$$

$$= \frac{1}{2} TNL_t^A - \frac{\beta}{2} TL_{t+\Delta t}^A + \left(TNL_t - \frac{1}{2} TNL_{t-\Delta t} - \frac{\beta}{2} TL_t \right)^D$$

$$T_{t+\Delta t}^A = -\frac{\beta \Delta t}{2} TL_{t+\Delta t}^A + Y$$

$$Y = \frac{\Delta t}{2} TNL_t^A + \left\{ T_t + \Delta t \left(TNL_t - \frac{1}{2} TNL_{t-\Delta t} - \frac{\beta}{2} TL_t \right) \right\}^D$$

Semi-Implicit Semi-Lagrange

Momentum equations :

$$U_{t+\Delta t}^A = \tilde{U} - \frac{\beta\Delta t}{2} UL_{t+\Delta t}^A$$

$$V_{t+\Delta t}^A = \tilde{V} - \frac{\beta\Delta t}{2} VL_{t+\Delta t}^A$$

$$\overline{VL} = (UL, VL) = \cos\theta \nabla (AT + R_d T_r \ln p_s), \quad q = \ln p_s$$

rescaling to (u, v) and taking the divergence :

$$D^+ = X - \frac{\beta\Delta t}{2} \nabla^2 (AT^+ + R_d T_r q^+)$$

$$\text{where} \quad X = \nabla \cdot \left[\frac{1}{\cos\theta} (\tilde{U}, \tilde{V}) \right]$$

Thermodynamic and surface pressure equations :

$$T_{t+\Delta t}^A = Y - \frac{\beta\Delta t}{2} \tau D_{t+\Delta t}^A$$

$$q_{t+\Delta t}^A = Z - \Delta t S \cdot D_{t+\Delta t}^A$$

$$\text{let } F^+ = F \frac{A}{t+\Delta t}, \quad b = \frac{\beta \Delta t}{2}, \quad r = R_d T_r$$

$$D^+ = X - b \nabla^2 \left[A (Y - b \tau D^+) + r (Z - \Delta t S \cdot D^+) \right]$$

$$\left[I + b \nabla^2 (-b A \tau - r \Delta t S) \right] D^+ = X - b \nabla^2 (A Y + r Z)$$

$$\left[I + \frac{n(n+1)}{a^2} b (b A \tau + r \Delta t S) \right] D^+$$

SI(1)

$$= X + \frac{n(n+1)}{a^2} b (A Y + r z)$$

$$\text{let } D_{-m} =$$

$$I + \frac{n(n+1)}{a^2} \left[\left(\frac{\beta \Delta t}{2} \right)^2 A \tau + \left(\frac{\beta \Delta t}{2} \right)^2 \frac{R_d}{P^r} T_r \cdot \Delta p^r \right]$$

$$D_{-m} = I + \frac{n(n+1)}{a^2} \frac{\beta^2 (\Delta t)^2}{4} \left[A \tau + \frac{R_d}{P^r} T_r \cdot \Delta p^r \right]$$

and :

$$T_r \cdot \Delta P^r = \begin{pmatrix} T_1^r \\ \vdots \\ T_{Levs}^r \end{pmatrix} \begin{pmatrix} \Delta p_1^r & \dots & \Delta p_{Levs}^r \end{pmatrix}$$

$$\text{Since } S = \frac{\beta}{2P_S^r} \begin{pmatrix} \Delta p_1^r & \dots & \Delta p_K^r \end{pmatrix}$$

We can use subroutine `get_cd_hyb` with $\frac{\Delta t}{2}$ to get D_m for slg.

This result is identical to the Eulerian forward step formulation.

Eq. SI(1) can now be solved for D^+ .

$$D^+ = D_m^{-1} \left[x + \frac{n(n+1)}{a^2} \frac{\beta \Delta t}{2} (AY + R_d T_r Z) \right]$$

Finally T^+ and q^+ can also be isolated.

mass adjustment for slg

$$\tilde{D}^+ = X + \tilde{X} - \frac{\beta}{2} \Delta t \nabla^2 (A \tilde{T}^+ + RT_r \tilde{q}^+)$$

$$D^+ = X - \frac{\beta}{2} \Delta t \nabla^2 (AT^+ + RT_r q^+)$$

subtract

$$\tilde{D}^+ - D^+ = \tilde{X} - \frac{\beta}{2} \Delta t \nabla^2 \{A(\tilde{T}^+ - T^+) + RT_r(\tilde{q}^+ - q^+)\}$$

$$\tilde{T}^+ = Y + \tilde{Y} - \frac{\beta \Delta t}{2} \tau \tilde{D}^+$$

$$T^+ = Y - \frac{\beta \Delta t}{2} \tau D^+$$

subtract

$$\tilde{T}^+ - T^+ = \tilde{Y} - \frac{\beta \Delta t}{2} \tau (\tilde{D}^+ - D^+)$$

$$\tilde{q}^+ = Z + \tilde{Z} - \frac{\beta}{2} \Delta t S \tilde{D}^+$$

$$q^+ = Z - \frac{\beta \Delta t}{2} S \cdot D^+ \quad S = \frac{1}{P_s^r} (\Delta P_1^r \dots \Delta P_{Levs}^r)$$

subtract

$$\tilde{q}^+ - q^+ = \tilde{Z} - \frac{\beta}{2} \Delta t S \cdot (\tilde{D}^+ - D^+)$$

$$\text{let } D^* = \tilde{D}^+ - D^+, \quad T^* = \tilde{T}^+ - T^+, \quad q^* = \tilde{q}^+ - q^+$$

then

$$D^* = \tilde{X} - \frac{\beta}{2} \Delta t \nabla^2 (AT^* + RT_r q^*)$$

$$T^* = \tilde{Y} - \frac{\beta}{2} \Delta t \tau D^*$$

$$q^* = \tilde{Z} - \frac{\beta}{2} \Delta t S \cdot D^*$$

These equations have the same form as the semi-implicit system for the solution of the dynamic part.

$$D^* = \tilde{X} - \frac{\beta}{2} \Delta t \nabla^2 \left[A \left(\tilde{Y} - \frac{\beta}{2} \Delta t \tau D^* \right) + RT_r \left(\tilde{Z} - \frac{\beta}{2} \Delta t S \cdot D^* \right) \right]$$

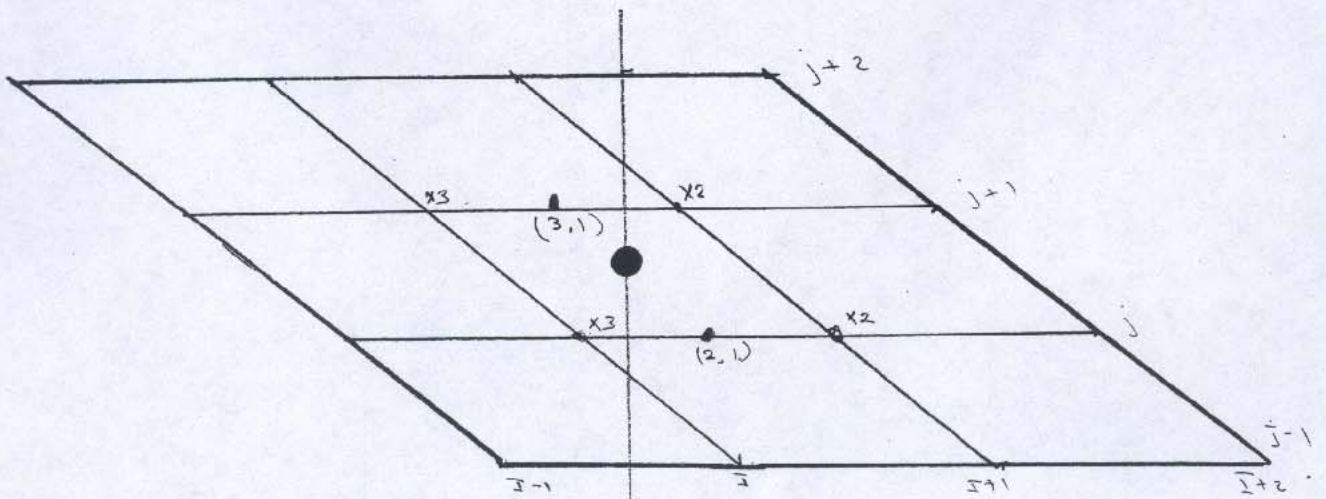
$$\left[I - \left(\frac{\beta}{2} \Delta t \right)^2 \nabla^2 \{ A \tau + RT_r S \cdot D \} \right] D^* = \tilde{X} - \beta \frac{\Delta t}{2} \nabla^2 [A \tilde{Y} + RT_r \tilde{Z}]$$

$$D^* = D_m^{-1} \left\{ \tilde{X} + \frac{n(n+1)}{a^2} \left(\beta \frac{\Delta t}{2} \right) [A \tilde{Y} + RT_r \tilde{Z}] \right\}$$

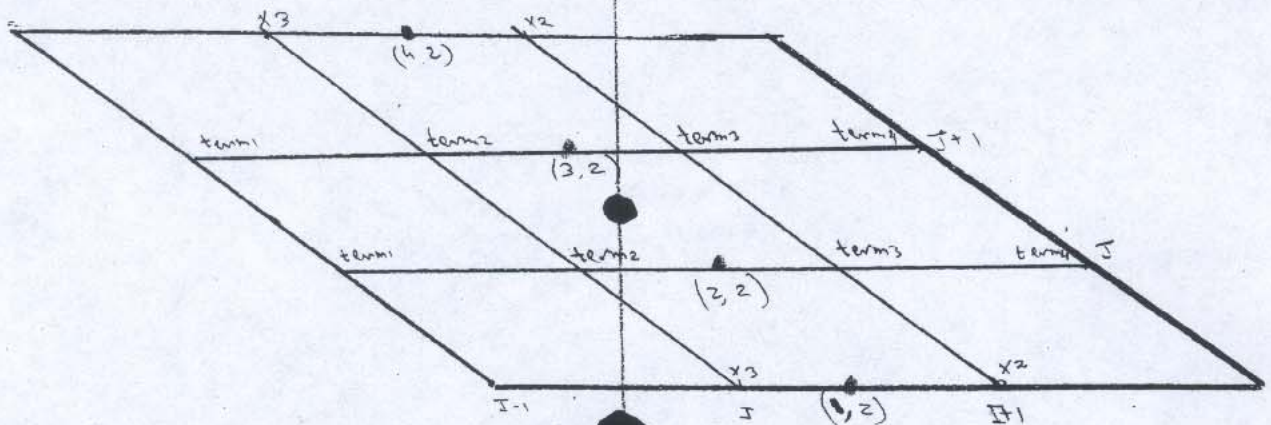
$$\tilde{D}^+ = D^* + D^+$$

$$\tilde{q}^+ = q^* + q^+$$

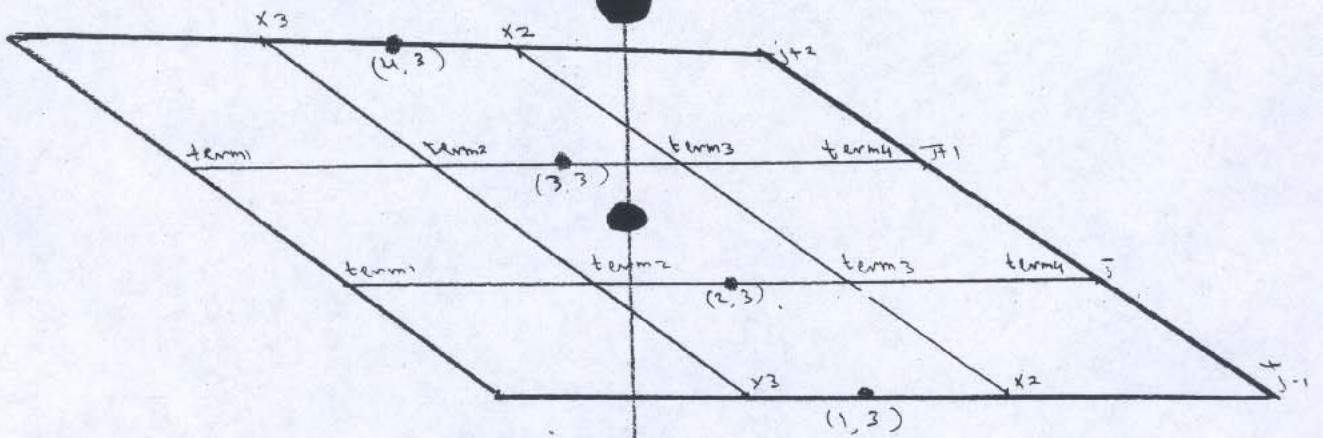
K-1



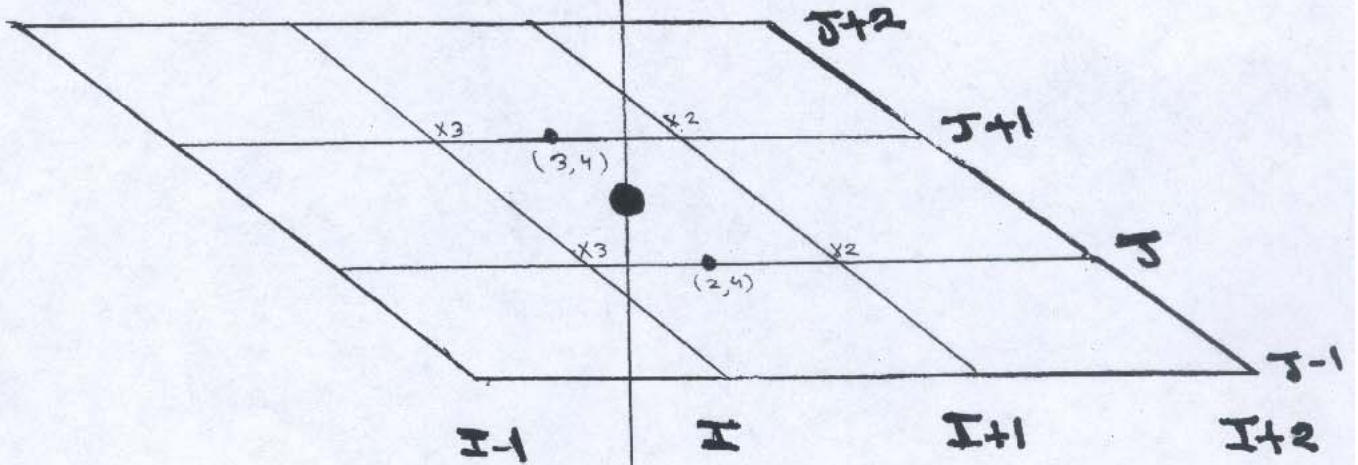
K



K+1

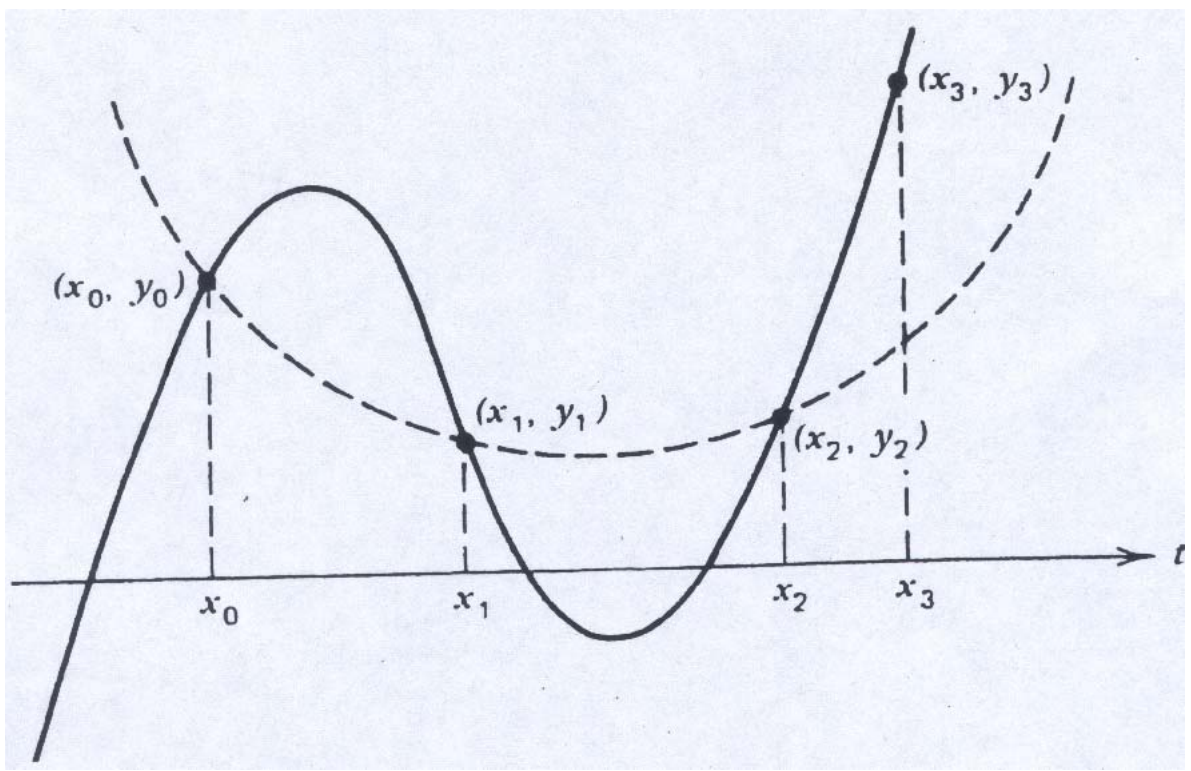


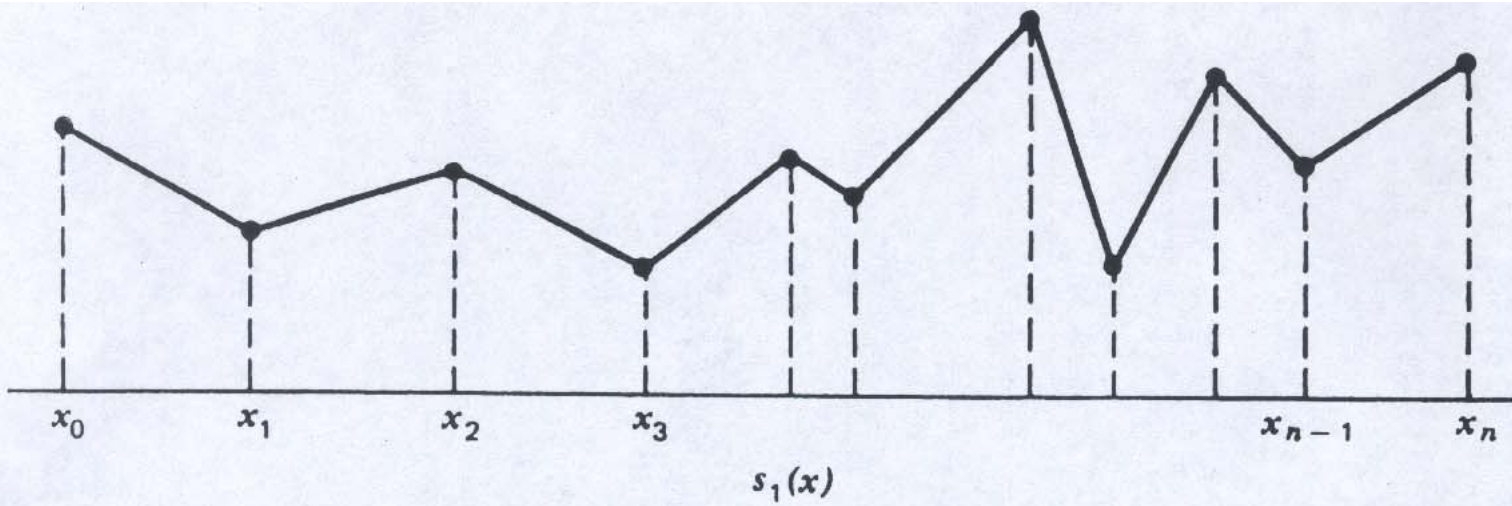
K+2



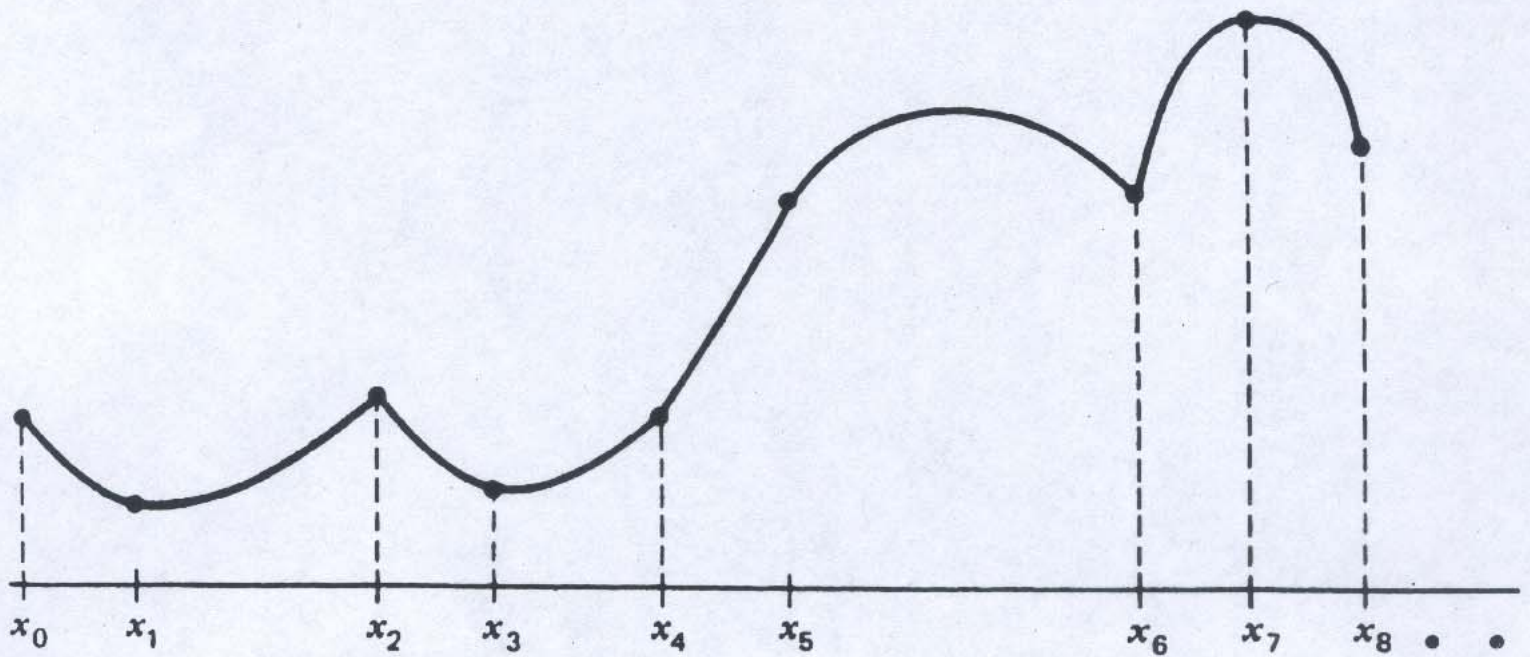
$$p_n(x) = \sum_{i=0}^n f(x_i) l_i(x),$$

$$l_i(x) = \frac{w(x)}{(x - x_i) w'(x_i)}$$
$$= \frac{(x - x_0) \cdots (x - x_{i-1})(x - x_{i+1}) \cdots (x - x_n)}{(x_i - x_0) \cdots (x_i - x_{i-1})(x_i - x_{i+1}) \cdots (x_i - x_n)}.$$





A piecewise linear Lagrange polynomial.



A piecewise quadratic Lagrange polynomial.

$$p(a) = f(a), \quad p'(a) = f'(a),$$

$$p(b) = f(b), \quad p'(b) = f'(b)$$

$$p(t) = a_0 + a_1t + a_2t^2 + a_3t^3$$

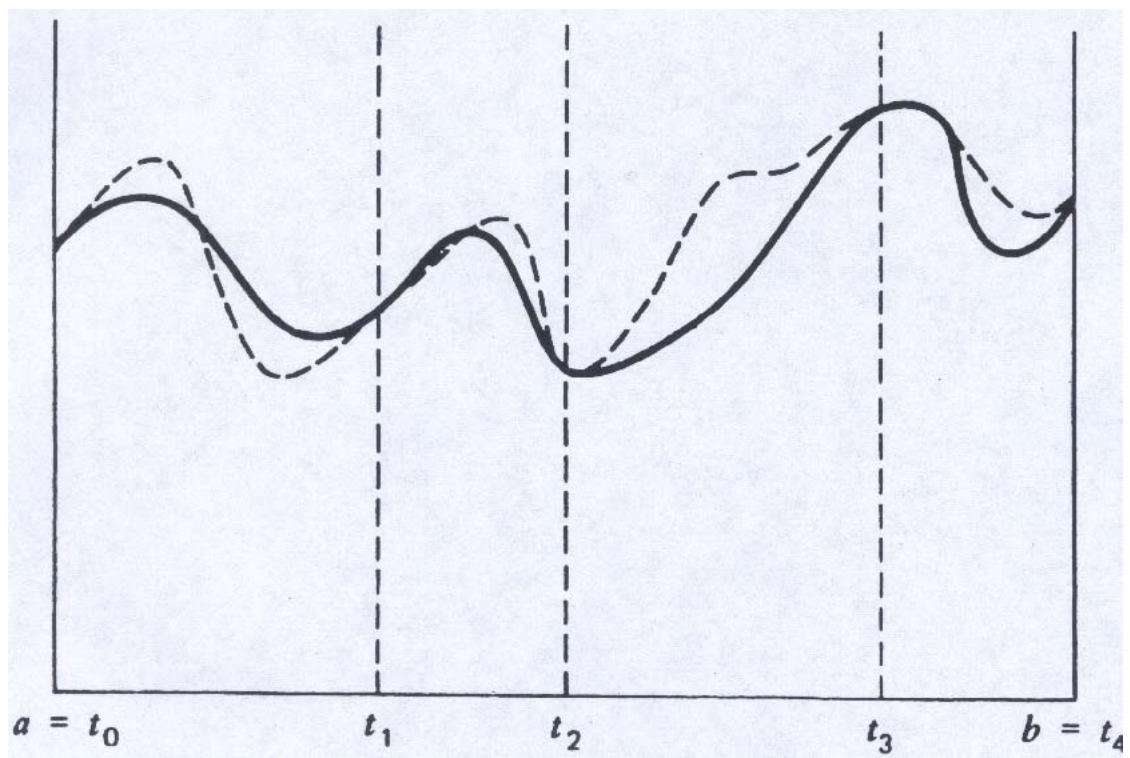
$$p(t) = p(a)\phi_1(t) + p(b)\phi_2(t) + p'(a)\psi_1(t) + p'(b)\psi_2(t),$$

$$\phi_1(t) = \frac{(t-b)^2[(a-b) + 2(a-t)]}{(a-b)^3}$$

$$\phi_2(t) = \frac{(t-a)^2[(b-a) + 2(b-t)]}{(a-b)^3}$$

$$\psi_1(t) = \frac{(t-a)(t-b)^2}{(a-b)^2}$$

$$\psi_2(t) = \frac{(t-a)^2(t-b)}{(a-b)^2}$$

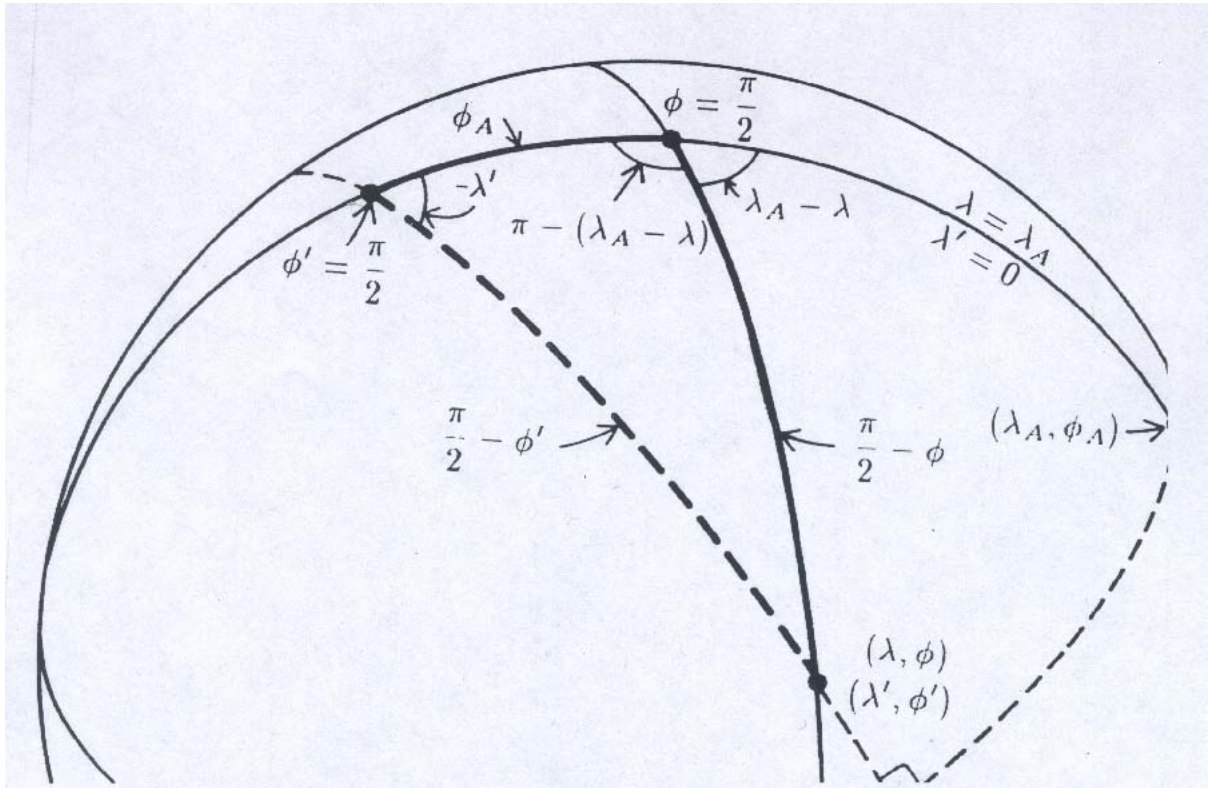


$$\lambda_D + \int_{(\lambda_D, \phi_D, t)}^{(\lambda_A, \phi_A, t + \Delta t)} \frac{u(\lambda, \phi, t)}{\cos \phi} dt = \lambda_A$$

$$\phi_D + \int_{(\lambda_D, \phi_D, t)}^{(\lambda_A, \phi_A, t + \Delta t)} v(\lambda, \phi, t) dt = \phi_A$$

$$\lambda_D^{k+1} = \lambda_A - \frac{1}{2} \Delta t \left[\frac{u(\lambda_A, \phi_A, t + \Delta t)}{\cos \phi_A} + \frac{u(\lambda_D^k, \phi_D^k, t)}{\cos \phi_D^k} \right]$$

$$\phi_D^{k+1} = \phi_A - \frac{1}{2} \Delta t [v(\lambda_A, \phi_A, t + \Delta t) + v(\lambda_D^k, \phi_D^k, t)]$$



$$\sin \phi' = \sin \phi \cos \phi_A - \cos \phi \sin \phi_A \cos(\lambda_A - \lambda)$$

$$\sin \lambda' \cos \phi' = -\sin(\lambda_A - \lambda) \cos \phi.$$

$$\lambda'_D = -\Delta t \left[\frac{u'(\lambda'_*, \phi'_*, t + \Delta t/2)}{\cos \phi'_*} \right]$$

$$\phi'_D = -\Delta t \left[v' \left(\lambda'_*, \phi'_*, t + \frac{\Delta t}{2} \right) \right]$$

$$\lambda'^{k+1}_* = -\frac{\Delta t}{2} \left[\frac{u'(\lambda'^k_*, \phi'^k_*, t + \Delta t/2)}{\cos \phi'^k_*} \right]$$

$$\phi'^{k+1}_* = -\frac{\Delta t}{2} \left[v' \left(\lambda'^k_*, \phi'^k_*, t + \frac{\Delta t}{2} \right) \right]$$

Performance

	T382k64	T510k64
Euler	1.	
Lagrange la_qb	3.	
Lagrange la_lb	3.75	1.95
Hermite la_qb	2.87	
Hermite la_lb	3.46	2.04